

Shelter Island II

Gell-Mann

~~Beta Renormalization~~ of present theory.

$W = X$ boson discovered.

Parameters, esp. ~~of~~ coupling of Higgs boson

Why 3 families? How many Higgs bosons? Unified Yang-Mills theory. Is renormalizability a good criterion?

Pecci-Quinn theory Higgslet, also called axis axion. Could have very small coupling, large $\langle \Phi \rangle$, then unobserved. "Dial" parameters to near zero. Limit of high energy.

10^{15} GeV. Coupling constant $\approx 1/40$

m_τ / m_b is OK

m_μ / m_s not OK, if our m_s is OK.

$SO(10)$ somewhat prettier. It takes guts to extrapolate GUTS to 10^{15} GeV. Is $10^{15} \approx$ Planck mass of 10^{19} ?

Proton decay, family transitions by neutral ~~currents~~ currents. Major $>$ antimatter to be explained. SU is

SO_{10} gives 3 neutrino masses.

~~Renorm QCD~~

Small masses compared to unification mass.

$$\frac{m}{M_{\text{unif}}} \sim \exp\left(-\frac{\text{const}}{\alpha_{\text{unif}}}\right)$$

Since $\alpha_{\text{unif}} \cong \frac{1}{40}$, this explained OK

If you dial ~~masses~~ _{constant} small, they stay small in renormal.

Fundamental fields to be called haptons, others composite. Higgs might be comp.

$n=1$ supersymmetry, works to some extent. Goldstino breaks by supersymmetry. Quantities can be only be estimate.

Supergravity $n=8$, many particles $1+8+28+70$ etc. particles, completely renormalizable.

Curiosities: Triplet of color & triplet of families

Ellis, Gaillard, Zumino (1980). All known particles are composite except graviton.

String theories No. of states, Regge trajectories

$N=4$ supergravity & $N=4$ Super-Yang-Mills

Super-strings II, no open strings. $N=8$ supergrav., no supermass. Appears finite.

Renormalization of 0-point energy. Takes revenge³
in gravity. Cosmological constant in astronomy 10^{-120} .

Quantum (bubbles) fluctuations (Hawking)

Inflationary Universe Linde-Guth.: Cosmological
constant jumps from finite to zero. Hawking: no problem.

Weinberg

Calc. of Fine Structure Constants

Charges from Extra Dimensions

Kaluza 21, Klein 26, Mandel

5 dimensions Minkowski $4 \times S^1$

S^1 is tight circle, radius ρ

$\hbar$ mass $1/\rho$ ρ undetermined

Charge $e \sqrt{16\pi G} / \rho$ for $l=1$

$1/\rho \approx 10^{17} \text{ GeV}$, $\rho \sim 10^{-31} \text{ cm}$

Higher Dimensions $4 + N$ dimensions

De Witte 1964, many other papers

Symmetry bringing you back to same pt.

rms of circumference replaces $2\pi\rho$

Some progress in calc. ratios of coupling constants.

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f massless matter fields in $4+N$ dimensions

Energy density $\propto \rho^{-4-N}$

$$\rho^2 \sim G f$$

$$e \approx \frac{\sqrt{G}}{\rho} \approx \frac{1}{\sqrt{f}}$$

Classical gravitat. fields, large f .

Gauge coupling constants < 1 .

Size of fine structure constant has not changed since quasars.

If N is odd, no possible additions to $\mathcal{L}$ can eliminate dimen

$$g^{vac} = \begin{pmatrix} \eta_{\mu\nu} & 0 \\ 0 & \tilde{g}_{mn}(y) \end{pmatrix}$$

$$\Gamma_{eff} = \int d^4x \mathcal{V}_{eff}(\tilde{g})$$

$$\mathcal{V}_{eff} = \int d^N y \sqrt{\tilde{g}} \frac{\tilde{R}(y) + \bar{\Lambda}}{16\pi G} + \mathcal{V}(\tilde{g})$$

matter loops

$\bar{\Lambda}$ cosmological constant in $4+N$ dimens.

$$\frac{N-1}{8\pi G \rho^2} = \mathcal{V}(\rho) - \frac{\rho}{N} \frac{d\mathcal{V}(\rho)}{d\rho}$$

$$\rho^2 = \frac{8\pi G (N+4) C_N}{N(N-1)}$$

C_N 's of order 1.

Quantize in space $\mathcal{Q}$. Integrals obtained

Appelquist & Chodos calc. results, for bosons $C_N > 0$ for $N=3$ to 19 , negative for $N=1$ or ≥ 21 .

For fermions and bosons, all OK if $N=7, 11$ etc.

Embarrassing that C_N is small, bec.

$$g \simeq \frac{N}{\sqrt{C_N}}$$

This is couplg. const. at 10^{17} GeV = $\frac{10^{19} \text{ Planck}}{\sqrt{16\pi}} g$

Curvature term $\propto \rho^{N-2}$

$\Lambda \propto \rho^N$

Matter term $V \sim r^{-4}$ prevents collapse

Increase temp. $\frac{3}{2}$ effex. : V gets term $\rho^N \mp N+4$, is < 0

Crit. temp. order 10^{17} GeV (Kaluza-Klein).

2-loop divergences, 1st approx. is good.

He thinks there is no finite th. incl. gravity.

$p\bar{p}$ machine $1.7 \cdot 10^9$

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$W \rightarrow e \nu$

Calorimeter

1+1 energy, large angle isolated electron

Transverse m. of electron 20-50 GeV

↳ of ν (missing mass)

about 20 events of W

Max. likelihood fit $M_W \approx 81 \pm 2$ GeV

Rapidity fits V-A

Transverse momentum of W abt. .

$\sigma_B = 2$ nb

Candidates for Z^0

1. One event $\mu^+ \mu^-$ $M_Z = 95 \pm 10$ GeV

2. $e^+ e^-$, one E = 49 GeV

Possibly $M_Z = 102 \pm 7$

3. $e^+ e^-$ $M = 93$

4. " 982

Wiki

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Kramers 1940 (Ann Arbor) suggested calc.

diff. betw. E (2s, H). Shelver Isl. gave actual Lamb shift.

Lamb Old exps. Lamb's history

1938-51 Columbia

H, D $n=2$

He⁺ $n=2$

QED K, L

1951-56 Stanford

H $n=3, 4$

He $2^3P_1, 3^3P_1$

1956-62 Oxford

Laser theory

1962-74 Yale

Very high Rydberg states of He

Experiments, Shift 1957. + MHz

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Trichwasser, D.L. .77(6)

R.S. (70) .90(6)

L.P. (75) .892(20)

N.A.V. (79) .862(20)

L. Pipkin (81) .845(9)

Calculation

Ericson (71) .930(10)

Mohr (75) .884(13), with $\tau_p = .862(12)$

Kinoshita: Mohr is correct

Kinoshita

Heard abt. $g = 2$ in Japan

1947 Nafe, Nelson, Rabi $hfs \neq$ Dirac

Breit suggested $g \neq 2$

Kusch, Foley after

1948 Schwinger $\alpha/2\pi$

1957 Sommerfeld 4 α order

Sumra $g_{\mu} - g_e$

1967 $\alpha_{\mu} - \alpha_e$ 6 α order, by renormaliz. group & mass singularity theorem

Parker, Taylor, Langenberg

1969-82 Kinoshita etc.

$$a_e^6 = 1.176 \ 5 \quad (\alpha/\pi)^3$$

$$= 24.48 \quad (\alpha/\pi)^3$$

72 Electron-weak

77 Van Dyck, Schwinger, Dehmelt, free electron

$$a = 1, 159, 652, 40 \ (20) \cdot 10^{-11}$$

81 a_μ at Com

$$81 \text{ Dehmelt} \quad 652 \ 200 \ (40) \cdot 10^{-12}$$

$$222 \ (50)$$

Theory to 8th order $\sim 0.8 \ (1.5)$

$$\alpha^{-1} = 137.035 \ 963 \ (15) \text{ Williams, Olson 1991}$$

$$a_e \text{ theory} = 1.159 \ 652 \ 460 \ (47 \text{ theory, } 127 \text{ error in } \alpha)$$

$$a_\mu \text{ Expt. } 1, 165, 937 \ (12) \cdot 10^{-9}$$

$$\text{Theor. } 1, 165, 920 \ (3.5) \cdot 10^{-9}$$

of which $65.5 \ (3.4)$ is hadron contrib., weak inter. $2 \cdot 10^{-9}$

Improve error in hadron contrib.

Determine α from $g-2$

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$$\alpha (g-2) = 137.035992 \text{ (5 exp.) (5 theor.)}$$

$$\text{muonium hfs} \quad 988 \text{ (21 exp.) (2.7 theor.)}$$

$$\text{Josephson} \quad 563(15) \text{ Williams, Olsen}$$

$$\text{Quantized Hall} \quad 968(23) \text{ } \frac{1}{2}$$

μ hfs

$$R_{\text{HQS}} = \frac{v B}{i} = \frac{h}{e^2 n}$$

Self-energy

$$W = - \frac{2e^2}{3\pi\hbar c^3} \int_0^K \hbar dk \sum \frac{|U_{m,k}|^2}{E_n - E_m + \hbar} \quad \text{My paper}$$

$$W_0 = \int_0^K \hbar dk \underbrace{(U^2)_{nn}}_{\hbar} \quad \text{free electron}$$

U has only diagonal elements

$$(U^2)_{nn} = \sum_k |U_{nk}|^2$$

$$\begin{aligned} \therefore W - W_0 &= + \int_0^K dk \sum (E_n - E_m) \frac{|U_{mn}|^2}{E_n - E_m + \hbar} \\ &= \frac{2}{m^2} \sum (E_n - E_m) |p_{mn}|^2 \ln \frac{K}{|E_n - E_m|} \end{aligned}$$

(principal value)

$$K = \hbar c^2$$

$$\begin{aligned} \sum (E_n - E_m) |p_{mn}|^2 &= -\hbar^2 \int \psi_m^* \nabla V \cdot \nabla \psi_n d\tau \\ &= \frac{1}{2} \hbar^2 \int \nabla^2 V \psi_n^2 d\tau = 2\pi\hbar^2 e^3 Z \psi_n^2(0) \end{aligned}$$

Must then calc.

$$\frac{\sum (E_n - E_m) |p_{mn}|^2 \ln(E_n - E_m)}{\sum (E_n - E_m) |p_{mn}|^2} = \ln(E_n - E_m)_{Av}$$

$$(E_n - E_m)_{Av} = 17.8 R_y$$

PR 72, 339 (47)

~~F. D. Lee~~ Nambu

Tomonaga started covariant theory in '43, was ready when Lamb shift was discovered

T. D. Lee

Time as dynamical variable

Relativ. Q Mech $\phi(r, t)$ operators, $\vec{r}, t$ parameters.

Discrete mechanics, 'classical

Within T, have N space-time points

$$A = \sum_n \left\{ \frac{1}{2} v_n^2 - \frac{1}{2} (V(r_n) + V(r_{n-1})) \right\} (t_n - t_{n-1})$$

$$\frac{\partial A}{\partial t_n} = \frac{1}{2} v_n^2 = \frac{1}{2} [V(\vec{r}_n) + V(\vec{r}_{n-1})] = E_{n+1}$$

$$v_n = \frac{\vec{r}_n - \vec{r}_{n-1}}{t_n - t_{n-1}}$$

Reproduce uniform motion, harmonic oscillator.

Non-Rel quantum

$$\mathcal{G}(T) = \int e^{iA} \prod_n d^3 r_n dt_n$$

$$A = \sum_n \left\{ \frac{(r_n - r_{n-1})^2}{2(t_n - t_{n-1})} - \frac{1}{2} (t_n - t_{n-1}) [V(r_n) + V(r_{n-1})] \right\}$$

$(\vec{r}_n)$ op (t_n) op

Unitarity is questioned. There may be problems only in order $e^{-T/l}$, where l is elementary time unit.

Lamb shift exp. requires that $l < 1.6 \cdot 10^{-14}$ cm

Relativistic QM

Massless scalar

$$G = e^{-A} \int \prod d^4 x_i d\phi_i$$

$g = g^{-4}$ fundamental constant

$$A = \frac{1}{2} \sum \lambda_{ij} (\phi_i - \phi_j)^2 - j_i \phi_i$$

Like Kirchhoff law

Draw a hypersphere thru any group of $D+1$ sites, then $D+1$ sites form a cluster.

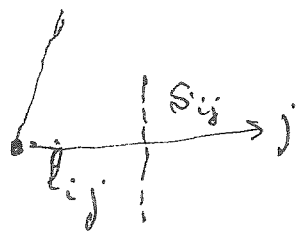
Theorem: Form a D -simplex for every cluster, using its sites as vertices. Non-overlapping simplices.

$$\sum \lambda_{ij} (\phi_i - \phi_j) = j_i \quad \sum \text{all sites linked to } i.$$

~~When $j_i = 0$, then $\phi_i = \phi_0 + \vec{E}$~~

$$\text{Choose } \lambda_{ij} = \frac{S_{ij}}{l_{ij}}$$

~~all~~ ϕ_i is contin. fn. of $\vec{r}_1, \dots, \vec{r}_N$



Gives no. of links, plaquettes

Gauge theory. Wilson uses cubic lattice, gives confinement easily. Lee's random lattice should give transition to weak coupling. not easily.

Calculated transition from strong to weak coupling, is smooth.

$SU(3)$, 4^4 lattice. Av. sweeps for each β is 40,000
 Gets very slow $\&$ max. in specific heat as fun. of β .

Jackiw

Non-perturbative calc. Skyrme suggestions.
 Infinities tamed only recently.

Local field dynamics, synthesis of QCD and flavor.

Stater need spontaneous symmetry breaking

Solitons: 1-D: kink, 2-D vortices, 3-D monopoles.
 Solve classical field,

Weak coupling g :

Mass = g^{-2} · charact. mass

Stability: $\propto$ energy barrier topologic g^4 no.

Interactions: Soliton-soliton g^{-1}

Excitation.

boson of integral spins $\leftrightarrow$ fermions

Fractional fermion no.

2. Processes

Yang-Mills th. tunneling betw. states of same energy.

Band formation - Tunneling

Learned topology from mathematicians

1. Yang-Mills gauge transform.

Non-Abelian gauge theory $U = e^{iT^a \theta_a}$

Not all deformable to unitary transf.

Large gauge transf. $|\text{state}\rangle \rightarrow e^{i\theta} |\text{state}\rangle$

Analog of Bloch theory $\psi(x+a) = e^{i\theta} \psi(x)$

Axial vector anomaly: conseq. of gauge th.

Physical Results

1. $\pi^0 \rightarrow 2\gamma$

2. quark - lepton balance

3. No chiral $U(1)$

Uncertain results:

4. CP violation by vacuum angle

5. Monopole

6. Mon.-induced baryon decay

7. Color breaking by monopoles

Domain boundaries are solitons in solid state physics

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Poly-acetylene 2 ground states



2 Solitons



Count links betw. solitons

Quantization

1. Dynam. variables : E , ang. momentum $\sim \hbar$
2. Soliton number : No $\hbar$, topology
3. Qu. of parameters in Lagrangian

Dirac monopole

$$\mathcal{L} = \frac{1}{2} m \dot{r}^2 + \frac{e}{c} \vec{A} \cdot \dot{\vec{r}} \cdot t$$

Action, $\mathcal{L}$ not invariant. Charge of $R = 2\pi \hbar$

$$\therefore eg = \hbar c \frac{n}{2}$$

(2) 3-D gauge theory

Eq. of motion gauge invariant

$$\text{Action} = \int d^3x \mathcal{L} \rightarrow \text{Ave.} + n \frac{8\pi^2}{g^2} w$$

$$\frac{4\pi m}{g^2} = kn$$

(3) N Weyl fermion doublets

Integrate out fermions

$SU(2)$ chiral currents.

Fermion no. conserved, but not gauge invariant.

(4) Effective current algebra

Wess-Zumino action, constant $\pm$

Physical results

2) Topological mass of gauge field.

4) Anomaly term in σ models

Skyrme: Non-linear σ -model

Witten, Balachandran find low levels of baryon

6' Hoofs

Elementary particles & black hole

Entropy of black hole $S = 4\pi M^2$ (perhaps divide by M)

Particle $\lambda_{\text{comp}} \sim 1/M$ Black hole $R \sim M$

Planck mass $\lambda \sim R$ (Schwarzschild)

Emission of γ 's by black hole

Marshall

~~the~~ Cosmic ray meson at sea level (37); J.G. Wilson found $\sigma_{\text{sat}} \lesssim 1\%$ of theory, σ_{prod} in upper atm. $\approx \sigma_{\text{me}}$ (Schwarz).
Discrepancy (Nordheim & Hebb, 1939). Conversi, Pancini & Piccioni (1945-7)

2. Nuclear force: $1/r^3$ divergence, Møller-Rosenfeld, Strong coupling theory (Pauli etc.)

3. Japanese (Yukawa, Tomonaga, Sakata, Inoue). Paper reached U.S. at end 47.

4. Fermi, Teller & Weisskopf found slowing down time 10^{-13} sec. Real discrepancy, Oppie instruction to Shelter Island $\pm$

5. Marshall's 2-meson theory. Powell π meson. Underground penetration. Kink at 1000 m due to π lifetime.

6. Clarification (1948)

Planck Mass

$$\frac{GM}{r^2} = \frac{v^2}{r}$$

$$v^2 = GM/r$$

$$v^2 r^3 = G M r^2$$

$$M v r = h$$

$$M^2 v^2 r^2 = h^2 = G M^3 r^3$$

Handing 3/6/83

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~~Super~~ $N=8$ Supergravity could be the ultimate theory

- 1) It may be finite to all orders
- 2) It unifies gravity w. ~~matter~~ other forces

But

1) Gd. state is anti-de Sitter space w. large negative
 cosmo, $\Lambda = -e^2 M_P^2$

2) Observed universe not supersymmetric

3) Not correct spectrum of particles

Break symmetry:

One-loop divergences

$$\Delta I = A \chi + B \int \Lambda^2 dV$$

χ = Euler number

$B = 0$ for $N \geq 5$

$A = 3 - N$ for $N \geq 3$

Add counterterm $-K\chi$ to effective action

Topological coupling constant K obeys renormalization group eqn; soln.:

$$K = K_0 - 5 \log l_x / l_P$$

l_P Planck length, l_x typical length scale.

Fluctuations of .

$$\Lambda = \Lambda_0 + \underbrace{\sqrt{g}}_{\text{fluctuations?}} + 8\pi \underbrace{\frac{\mu_0}{M_P^2}}_{\text{bubbles}}$$

- 1) Either form or scalar field w/o kinetic term leads to a range of different values of $\Lambda_{\text{apparent}}$ in path integral
- 2) Negative values of $\Lambda_{\text{apparent}}$ are suppressed
- 3) Positive values lead to a small Euclidean 4-volume
- 4) Most favored is $\Lambda = 0$.

Guth Linde Inflationary Universe

1. One or more phase transitions

$$T = 10^{14} \text{ GeV}$$

2. Magnetic monopoles exist

$$M = 10^{16} \text{ GeV}$$

3. Baryon finite life $10^{30} - 10^{33} \text{ ys.}$

Standard theory

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi}{3} G \rho, \quad \frac{\dot{R}}{R} = H$$

Metric increasingly flat as $t \rightarrow 0$

Mass density dominated by radiation of

$$t^2 = \frac{1}{40} \frac{M_p}{t}$$

$$\rho = 10^{75} \text{ g cm}^{-3}$$

$$R = 10 \text{ cm}$$

Problems 1. Magn. monopoles, topological knots

$$n_M = \xi^{-3} \quad \xi = \text{correlation length}$$

$$\xi <$$

$n_M/s \geq 10^{-13}$, $s = \text{entropy density}$
remains constant since then, gives

$$\Omega > 10^{12}$$

2. Horizon problem. Large-scale homogeneity.

* Rad. decoupled at $T \simeq 4000^\circ \text{K}$. when $t \simeq 10^5 \text{ yr}$.

Horizon length = distance of light travel since $t=0$
= 90 times t at decoupling

3. Flatness Problem

Today, $0.01 < \Omega < 10$

$$\rho_c = \frac{3H^2}{8\pi G}$$

Then at $T = 10^{14} \text{ GeV}$ must have

$$\Omega \approx 1 \pm 8(10^{-49})$$

bec. Ω gets increasingly far from 1

4. Smoothness Problem

Early universe very unstable vs. perturbations.

$$\frac{\delta \rho}{\rho} \approx \frac{1}{\sqrt{N}} \sim 10^{-29} \text{ originally}$$

$\frac{\delta \rho}{\rho}$ prop. to t , grows to 10^{+15}

If you start at 10^{-25} s at $T = 10^9$ GeV

New Inflationary Universe

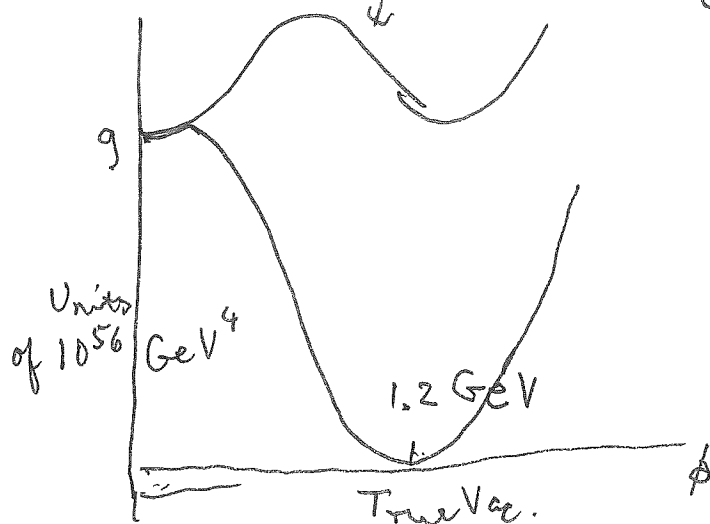
A single bubble can (probably) expand enough to encompass

SU_8 Coleman-Weinberg Potential

$$\underline{\Phi} = \sqrt{\frac{2}{15}} \phi \begin{pmatrix} 1 & & 0 \\ & 1 & \\ & & 1 \\ 0 & & -3/2 \end{pmatrix}$$

At origin, $\frac{\partial^2 V}{\partial \phi^2} = \frac{\partial^3 V}{\partial \phi^2} = 0$. Here have false vacuum
 $V(0) \approx 10^{56} \text{ GeV}^4$

At high temp. ↓



Crit. temp. when 2 minima equal. $\langle \Phi \rangle = 0$

Supercools w. $\langle \Phi \rangle = 0$ for $T < T_c$

$\rho \rightarrow \rho_0$ (energy density of false vac.)

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Flat Space (de Sitter).

All perturbations die out.

$$\dot{R}/R = \text{const.}$$

Observed universe

Baryon no. 10^{78-79}

Matter energy energy = 10^{78-80} GeV

Electric chge. = color chge. = ang. momentum ≈ 0

False vacuum is the source of all matter energy, and
Entropy

Universe devoid of all conserved quantities

Ref. Tryon; Brant, Englert & Spindel; Gott; Zurek

Phase Transition

ϕ fluctuates, $\xi \sim \chi^{-1} \sim 10^{-11}$ proton diameters

ϕ rolls classically down potential

$$\ddot{\phi} + 3 \frac{\dot{R}}{R} \dot{\phi} = - \frac{\partial V}{\partial \phi}$$

Exponential expansion much faster than rolling

$R(t)$ grows by factor $> 10^{25}$, ξ grows to $> 10 \text{ cm}$

Solution of problems

1. Monopole

2. Horizon: Univ. comes from a region of correlated ϕ . 23

3. Flatness

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi}{3} G \rho - \frac{k}{R^2}, \quad k \approx 1$$

While ϕ starts to roll, $\rho \approx \rho_0$, $k/R^2 \rightarrow 0$ by inflation

4. Smoothness

Qu. fluctuations in $\phi(\vec{x})$

$\frac{\delta \phi}{\rho}$ are scale-invariant

Mixing SU_5 gives 10^2 (too much)

Hawking & Moss ~ 1

Geometric hierarchy $\sim 10^{-4}$ (perfect)

Linde

Problems

1. SU_5 Coleman-Weinberg

2. Supersymmetric version

No baryons.

3. Supergravity coupled to matter (primordial inflation)

Inflation large, $G\rho/\rho \sim 10^{-4}$

Ellis, Nanopoulos

Monopoles? Reheating & baryons??

Monopoles produced after inflation

No phase transitions.

4. Improve model (Linde)

$$V = 3\mu^6 \left(1 - \alpha^2 \varphi^2 + \frac{\alpha}{4} \varphi^4 \right)$$

Duration of inflation $(3 \cdot 10^{11} \text{ GeV})^{-1}$

Extent of inflation e^{500}

Frie Pathology

$$\langle \varphi^2 \rangle = \frac{H^2}{4\pi^2} t$$

$H = \text{Hubble}$

After inflation,

$$\varphi = \langle \varphi^2 \rangle \sim H$$

Universe splits into many regions, domains, etc.

No monopoles

Chaotic Inflation

$$V(\varphi) = \frac{\lambda}{4} \varphi^4$$

$$\varphi = \varphi_0 \exp\left(-\frac{\sqrt{\lambda}}{6\pi} M_P t\right)$$

if $\varphi \geq M_P$

$$H^2 \geq 20 M^2(V)$$

$$V(\varphi) \leq M_{\text{Planck}}^4$$

4. Inflation inevitable in chaotic

Chaos $\rightarrow$ Universe w. small fluctuation

Linde Feynman QCD

1-Flavor quarks

2+1 dimensions (renormalization)

$SU(2)$

No dynamic quarks

Outline

Why gap to lowest glueball

Confinement of static quarks

Nature of string or bag

Color correlation hard to maintain

QED fluctuations

2-D space

3-D space

$$E \sim 1/L$$

$$B^2 L \sim 1/L$$

$$B^2 L^3 \sim 1/L$$

$$B \sim L^{-3/2}$$

$$B \sim 1/L$$

$$A \sim L^{-1/2}$$

$$A \sim 1/L$$

$$\int A \, dL = \text{Flux} \sim L^{1/2}$$

$$\text{Flux} \sim \ln L$$

$S(A)$ $\vec{A}$ is color vector.

Euclidean space-time

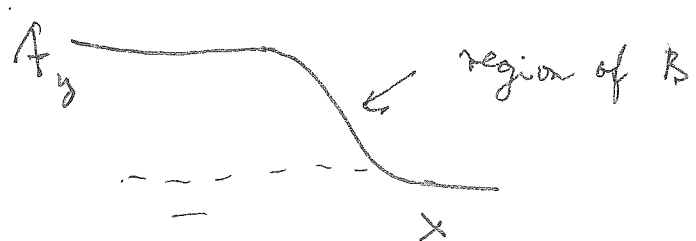
$$\text{Action } \dot{A}_x^2 + \dot{A}_y^2 + (\partial_x A_y - \partial_y A_x - A_x A_y)^2 = \mathcal{L}$$

$$\exp \int \mathcal{L} \, d\tau$$

If there is no state which ~~can decay~~ ^{can} decay.

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Why can A_y not diffuse? Bec. then $\partial_x A_y$ becomes large



$$A = \lambda \nabla \lambda$$